

The isentropic Euler system admits some plane wave superpositions

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1. Main Result

These velocities add, and pressures combine nonlinearly.

$$A_1 = \begin{bmatrix} u_1 & 0 & 0 & aw \\ 0 & u_1 & 0 & 0 \\ 0 & 0 & u_1 & 0 \\ aw & 0 & 0 & u_1 \end{bmatrix}, A_2 = \begin{bmatrix} u_2 & 0 & 0 & 0 \\ 0 & u_2 & 0 & aw \\ 0 & 0 & u_2 & 0 \\ 0 & aw & 0 & u_2 \end{bmatrix}, A_3 = \begin{bmatrix} u_3 & 0 & 0 & 0 \\ 0 & u_3 & 0 & 0 \\ 0 & 0 & u_3 & aw \\ 0 & 0 & aw & u_3 \end{bmatrix}$$

Isentropic Euler equations in symmetric [3] form

$$q_t + A_1 q_{x_1} + A_2 q_{x_2} + A_3 q_{x_3} = 0$$

where $q = (u_1, u_2, u_3, w)$ is a vector of velocities together with $w = a^{-1} \sqrt{\gamma p / \rho}$, proportional to the sound speed.

Here $p = k \rho^\gamma$ with k constant, $1 < \gamma < 3$, and $a = \frac{\gamma-1}{2}$.

THEOREM

There are solutions $q(x, t) = \sum_{j=1}^N f_j(x \cdot v_j, t) \begin{bmatrix} v_j \\ 1 \end{bmatrix}$
i.e.

$$u(x, t) = \sum_{j=1}^N f_j(x \cdot v_j, t) v_j,$$

and

$$\rho = \left(\frac{a}{\sqrt{k\gamma}} \sum_{j=1}^N f_j(x \cdot v_j, t) \right)^{\frac{1}{a}}.$$

Each $f_j(s, t)$ is a smooth inviscid Burgers solution

$$f_t + (1+a) f f_s = 0, \quad s \in \mathbb{R}, \quad 0 \leq t < T$$

and v_j are constant unit vectors with

$$v_i \cdot v_j = -a, \quad i \neq j.$$

The solutions exist for as long as the f_j remain smooth and $\sum f_j > 0$.

\mathbb{R}^d	$1 < \gamma < \frac{5}{3}$	$\gamma = \frac{5}{3}$	$\frac{5}{3} < \gamma < 2$	$\gamma = 2$	$2 < \gamma < 3$
\mathbb{R}^2	2	2	2	3	2
\mathbb{R}^3	3	4	3	3	2

The table shows the number N of vectors v_k available for various d and γ .

For details see

<http://www.math.cornell.edu/~bterrell/eulersuper.pdf>

2. Special cases

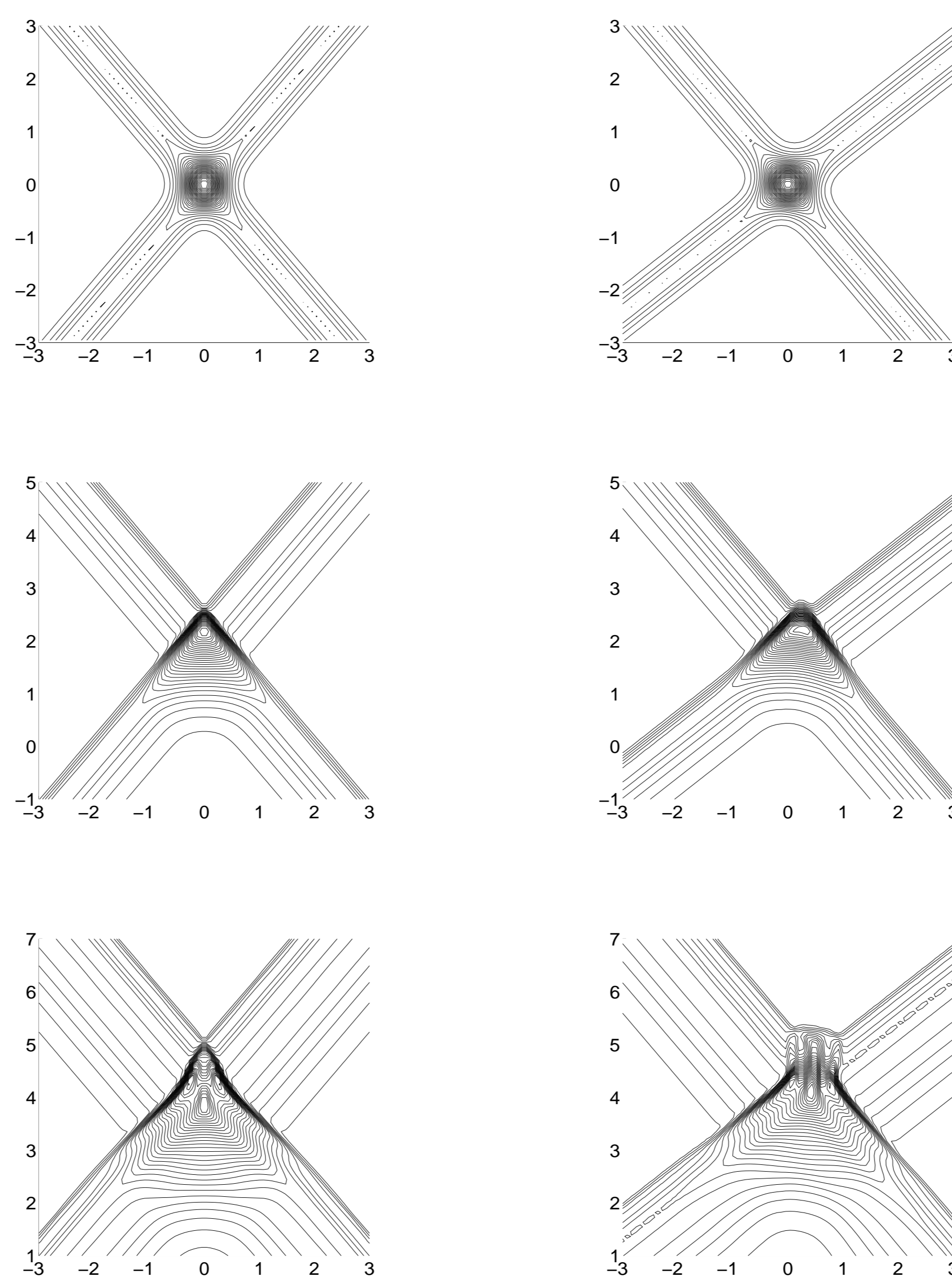
1) $\gamma = 2$, *shallow water model*, three vectors v_k coplanar at 120 degrees

2) $\gamma = \frac{5}{3}$, *monatomic gas*, four vectors v_k having the symmetry of a regular tetrahedron.

3. A computation

Illustration using a `clawpack`[2], [4] calculation. Density contours shown, with time increasing downward.

$\gamma = 1.4$, $k = 20000$, $f_1(s, 0) = f_2(s, 0) = 2 + 3 \exp(-5s^2)$, $f_3 = 0$. The solutions remain smooth almost until the time of the second figures.



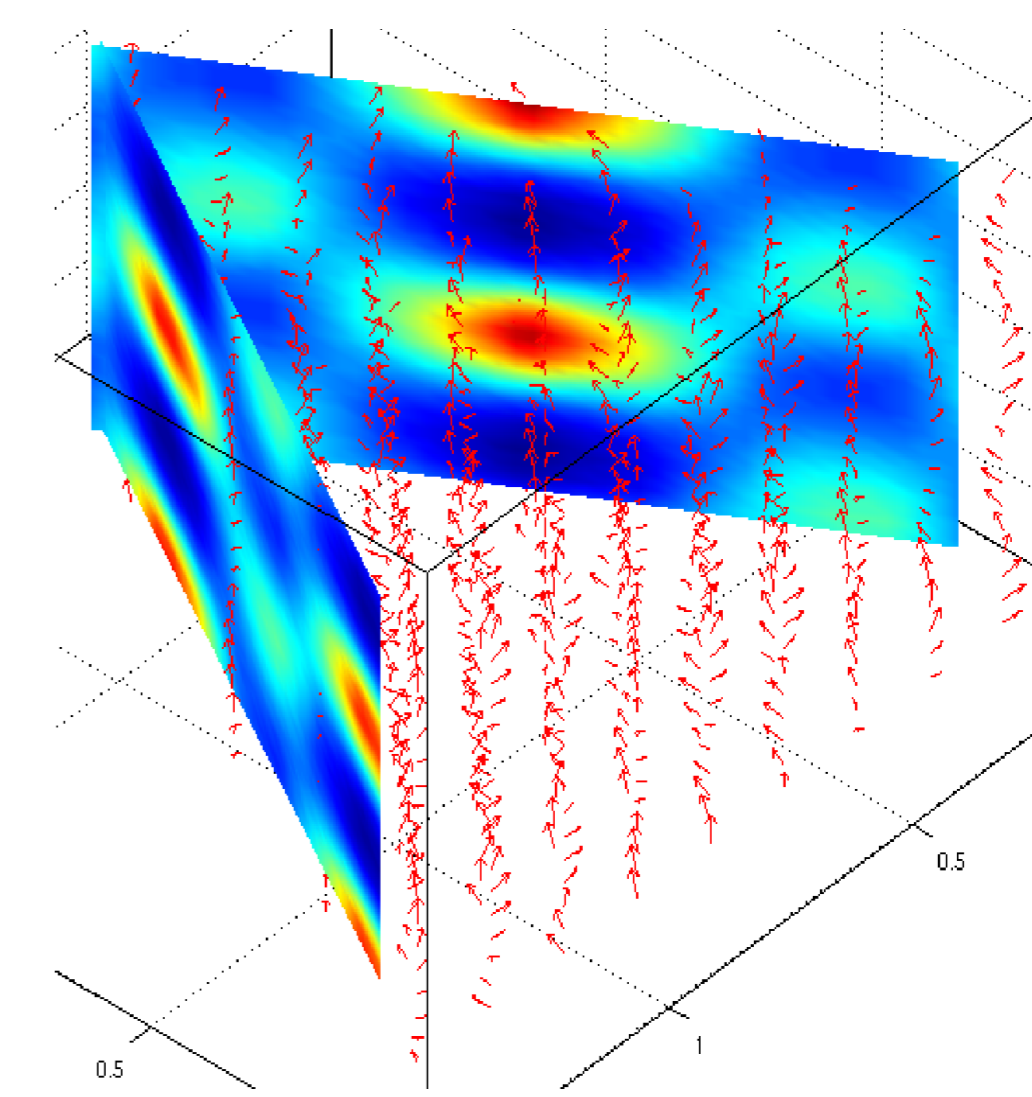
Left: angle is "correct".
Right: angle is "incorrectly" orthogonal.

4. Example in three dimensions

A flow in a triangular duct, tangential to the boundary.

$v_1 = (\sqrt{3}, 1, 1)/\sqrt{5}$, $v_2 = (-\sqrt{3}, 1, 1)/\sqrt{5}$, $v_3 = (0, -2, 1)/\sqrt{5}$, $\gamma = 1.4$

$f_1(s, 0) = f_2(s, 0) = f_3(s, 0) = f(s)$ has period p , and the duct sides have width $\sqrt{5}p/3$. Here $f(s) = 2 + 0.5 \cos(3\pi s)$.



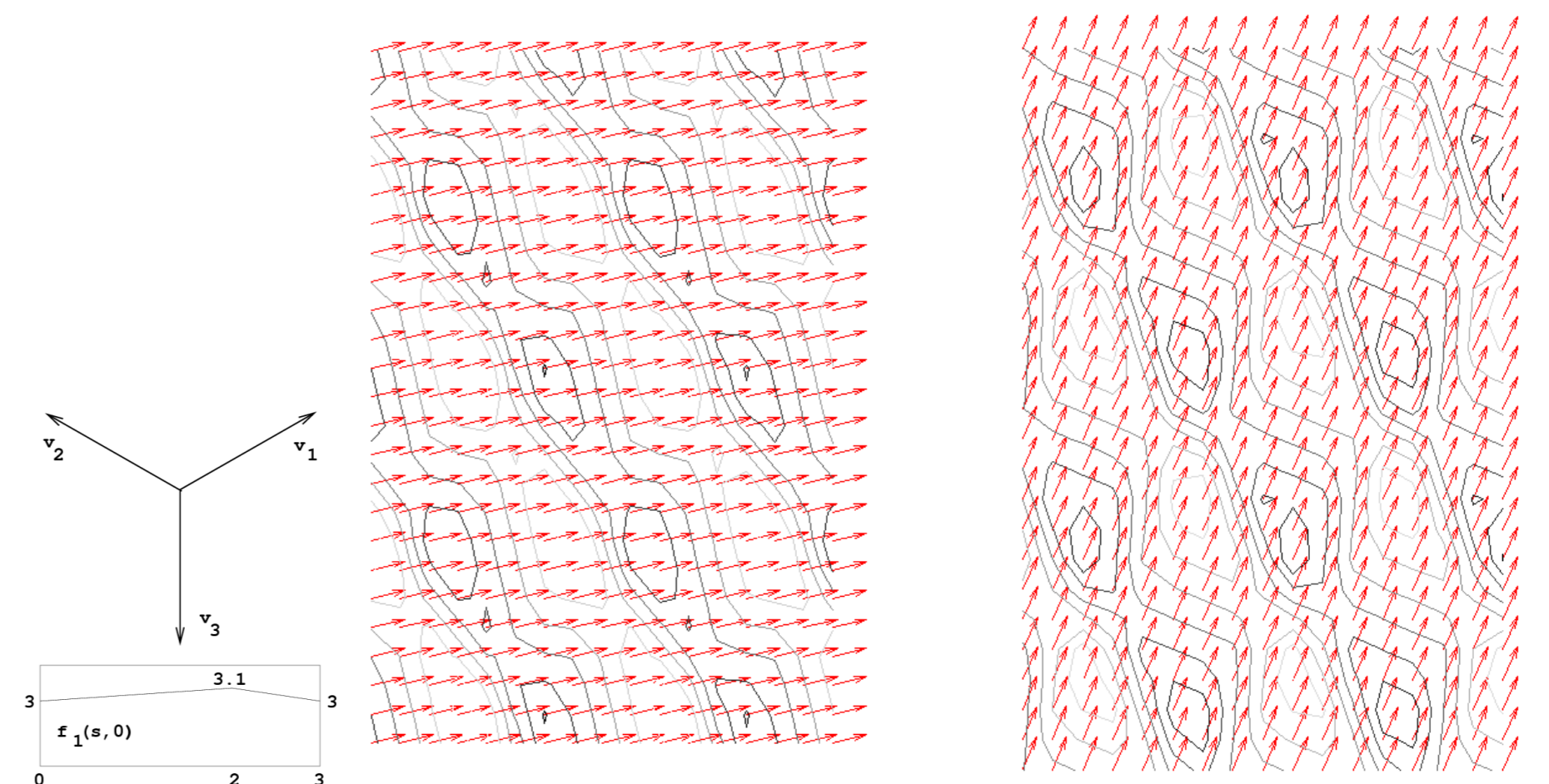
5. Example in two dimensions

An exact "shallow water" model, $\gamma = 2$, $g = 9.8$, $k = \frac{g}{2}$.

Left: $f_1(s, 0)$ is graphed, period 3, and the v_j indicated. The $f_3 = 0$.

Center: $f_2(s, t) = -0.4 f_1(s, -0.4t)$, $t = 0.5$, depth ~ 0.09 , speed ~ 3.8 .

Right: $f_2(s, t) = 0.6 f_1(s, 0.6t)$, $t = 0.5$, depth ~ 0.6 , speed ~ 2.7 .



References

- [1] J. M. BURGERS, *A mathematical model illustrating the theory of turbulence*, Adv. Appl. Mech., 1 (1948) pp. 171-179
- [2] <http://www.amath.washington.edu/~claw/>
- [3] P. D. LAX, *Hyperbolic Partial Differential Equations*, Providence, RI, American Mathematical Society, 2006.
- [4] R. J. LEVEQUE, *Finite Volume Methods for Hyperbolic Problems*, Cambridge University Press, 2002.